



**HAL**  
open science

# On Two Ways to Look for Mutually Unbiased Bases

Maurice Robert Kibler

► **To cite this version:**

Maurice Robert Kibler. On Two Ways to Look for Mutually Unbiased Bases. Acta Polytechnica, 2014, 54, pp.124-126. 10.14311/AP.2014.54.0124 . in2p3-00932804v2

**HAL Id: in2p3-00932804**

**<https://in2p3.hal.science/in2p3-00932804v2>**

Submitted on 18 Jan 2014

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.

## ON TWO WAYS TO LOOK FOR MUTUALLY UNBIASED BASES

MAURICE R. KIBLER\*

*Université de Lyon, Université Claude Bernard Lyon 1 et CNRS/IN2P3, Institut de Physique Nucléaire, 4 rue Enrico Fermi, 69622 Villeurbanne, France*\* corresponding author: [kibler@ipnl.in2p3.fr](mailto:kibler@ipnl.in2p3.fr)

**ABSTRACT.** Two equivalent ways of looking for mutually unbiased bases are discussed in this note. The passage from the search for  $d+1$  mutually unbiased bases in  $\mathbb{C}^d$  to the search for  $d(d+1)$  vectors in  $\mathbb{C}^{d^2}$  satisfying constraint relations is clarified. Symmetric informationally complete positive-operator-valued measures are briefly discussed in a similar vein.

**KEYWORDS:** finite-dimensional quantum mechanics, quantum information, MUBs, SIC POVMs, equiangular lines, equiangular vectors.

## 1. INTRODUCTION

The concept of mutually unbiased bases (MUBs) plays an important role in finite-dimensional quantum mechanics and quantum information (for more details, see [1–4] and references therein). Let us recall that two orthonormal bases  $\{|a\alpha\rangle : \alpha = 0, 1, \dots, d-1\}$  and  $\{|b\beta\rangle : \beta = 0, 1, \dots, d-1\}$  in the  $d$ -dimensional Hilbert space  $\mathbb{C}^d$  (endowed with an inner product denoted as  $\langle | \rangle$ ) are said to be unbiased if the modulus of the inner product  $\langle a\alpha | b\beta \rangle$  of any vector  $|b\beta\rangle$  with any vector  $|a\alpha\rangle$  is equal to  $1/\sqrt{d}$ . It is known that the maximum number of MUBs in  $\mathbb{C}^d$  is  $d+1$  and that this number is reached when  $d$  is a power of a prime integer. In the case where  $d$  is not a prime integer, it is not known if one can construct  $d+1$  MUBs (see [4] for a review).

In a recent paper [5], it was discussed how the search for  $d+1$  mutually unbiased bases in  $\mathbb{C}^d$  can be approached via the search for  $d(d+1)$  vectors in  $\mathbb{C}^{d^2}$  satisfying constraint relations. It is the main aim of this note to make the results in [5] more precise and to show that the two approaches (looking for  $d+1$  MUBs in  $\mathbb{C}^d$  or for  $d(d+1)$  vectors in  $\mathbb{C}^{d^2}$ ) are entirely equivalent. The central results are presented in Sections 2 and 3. In Section 4, parallel developments for the search of a symmetric informationally complete positive-operator-valued measure (SIC POVM) are considered in the framework of similar approaches. Some concluding remarks are given in the last section.

## 2. THE TWO APPROACHES

It was shown in [5] how the problem of finding  $d+1$  MUBs in  $\mathbb{C}^d$ , i.e.,  $d+1$  bases

$$B_a = \{|a\alpha\rangle : \alpha = 0, 1, \dots, d-1\} \quad (1)$$

satisfying

$$|\langle a\alpha | b\beta \rangle| = \delta_{\alpha,\beta} \delta_{a,b} + \frac{1}{\sqrt{d}}(1 - \delta_{a,b}) \quad (2)$$

can be transformed in the problem of finding  $d(d+1)$  vectors  $\mathbf{w}(a\alpha)$  in  $\mathbb{C}^{d^2}$ , of components  $w_{pq}(a\alpha)$ , satisfying

$$\overline{w_{pq}(a\alpha)} = w_{qp}(a\alpha), \quad p, q \in \mathbb{Z}/d\mathbb{Z} \quad (3)$$

$$\sum_{p=0}^{d-1} w_{pp}(a\alpha) = 1 \quad (4)$$

and

$$\sum_{p=0}^{d-1} \sum_{q=0}^{d-1} \overline{w_{pq}(a\alpha)} w_{pq}(b\beta) = \delta_{\alpha,\beta} \delta_{a,b} + \frac{1}{d}(1 - \delta_{a,b}) \quad (5)$$

with  $a, b = 0, 1, \dots, d$  and  $\alpha, \beta = 0, 1, \dots, d-1$  in (1)–(5). (In this paper, the bar denotes complex conjugation.) This result was described by Proposition 1 in [5]. In fact, the equivalence of the two approaches (in  $\mathbb{C}^d$  and  $\mathbb{C}^{d^2}$ ) requires that each component  $w_{pq}(a\alpha)$  be factorized as

$$w_{pq}(a\alpha) = \omega_p(a\alpha) \overline{\omega_q(a\alpha)} \quad (6)$$

for  $a = 0, 1, \dots, d$  and  $\alpha = 0, 1, \dots, d-1$ , a condition satisfied by the example given in [5]. The factorization of  $w_{pq}(a\alpha)$  follows from the fact that the operator  $\Pi_{a\alpha}$  defined in [5] is a projection operator.

The introduction of (6) in (3), (4) and (5) leads to some simplifications. First, (6) implies the hermiticity condition (3). Second, by introducing (6) into (4) and (5), we obtain

$$\sum_{p=0}^{d-1} |\omega_p(a\alpha)|^2 = 1 \quad (7)$$

and

$$\left| \sum_{p=0}^{d-1} \overline{\omega_p(a\alpha)} \omega_p(b\beta) \right|^2 = \delta_{\alpha,\beta} \delta_{a,b} + \frac{1}{d}(1 - \delta_{a,b}) \quad (8)$$

respectively. It is clear that (7) follows from (8) with  $a = b$  and  $\alpha = \beta$ . Therefore, (3) and (7) are redundant in view of (5) and (6). As a consequence,

Proposition 1 in [5] can be precised and reformulated in the following way.

**Proposition 1.** *For  $d \geq 2$ , finding  $d + 1$  MUBs in  $\mathbb{C}^d$  (if they exist) is equivalent to finding  $d(d + 1)$  vectors  $\mathbf{w}(a\alpha)$  in  $\mathbb{C}^{d^2}$ , of components  $w_{pq}(a\alpha)$  such that*

$$\sum_{p=0}^{d-1} \sum_{q=0}^{d-1} \overline{w_{pq}(a\alpha)} w_{pq}(b\beta) = \delta_{\alpha,\beta} \delta_{a,b} + \frac{1}{d}(1 - \delta_{a,b}) \quad (9)$$

and

$$w_{pq}(a\alpha) = \omega_p(a\alpha) \overline{\omega_q(a\alpha)}, \quad p, q \in \mathbb{Z}/d\mathbb{Z} \quad (10)$$

where  $a, b = 0, 1, \dots, d$  and  $\alpha, \beta = 0, 1, \dots, d - 1$ .

This result can be transcribed in matrix form. Therefore, we have the following proposition.

**Proposition 2.** *For  $d \geq 2$ , finding  $d + 1$  MUBs in  $\mathbb{C}^d$  (if they exist) is equivalent to finding  $d(d + 1)$  matrices  $M_{a\alpha}$  of dimension  $d$ , with elements*

$$(M_{a\alpha})_{pq} = \omega_p(a\alpha) \overline{\omega_q(a\alpha)}, \quad p, q \in \mathbb{Z}/d\mathbb{Z} \quad (11)$$

and satisfying the trace relations

$$\text{Tr}(M_{a\alpha} M_{b\beta}) = \delta_{\alpha,\beta} \delta_{a,b} + \frac{1}{d}(1 - \delta_{a,b}) \quad (12)$$

where  $a, b = 0, 1, \dots, d$  and  $\alpha, \beta = 0, 1, \dots, d - 1$ .

### 3. EQUIVALENCE

Suppose that we have a complete set  $\{B_a : a = 0, 1, \dots, d\}$  of  $d + 1$  MUBs in  $\mathbb{C}^d$ , i.e.,  $d(d + 1)$  vectors  $|a\alpha\rangle$  satisfying (2), then we can find  $d(d + 1)$  vectors  $\mathbf{w}(a\alpha)$  in  $\mathbb{C}^{d^2}$ , of components  $w_{pq}(a\alpha)$ , satisfying (9) and (10). This can be achieved by introducing the projection operators

$$\Pi_{a\alpha} = |a\alpha\rangle\langle a\alpha| \quad (13)$$

where  $a = 0, 1, \dots, d$  and  $\alpha = 0, 1, \dots, d - 1$ . In fact, it is sufficient to develop  $\Pi_{a\alpha}$  in terms of the  $E_{pq}$  generators of the  $\text{GL}(d, \mathbb{C})$  complex Lie group; the coefficients of the development are nothing but the  $w_{pq}(a\alpha)$  complex numbers satisfying (9) and (10), see [5] for more precisions.

Reciprocally, should we find  $d(d + 1)$  vectors  $\mathbf{w}(a\alpha)$  in  $\mathbb{C}^{d^2}$ , of components  $w_{pq}(a\alpha)$ , satisfying (9) and (10), then we could construct  $d(d + 1)$  vectors  $|a\alpha\rangle$  satisfying (2). This can be done by means of a diagonalization procedure of the matrices

$$M_{a\alpha} = \sum_{p=0}^{d-1} \sum_{q=0}^{d-1} w_{pq}(a\alpha) E_{pq} \quad (14)$$

where  $a = 0, 1, \dots, d$  and  $\alpha = 0, 1, \dots, d - 1$ . An alternative and more simple way to obtain the  $|a\alpha\rangle$  vectors from the  $\mathbf{w}(a\alpha)$  vectors is as follows. Equation (8) leads to

$$\left| \sum_{p=0}^{d-1} \overline{\omega_p(a\alpha)} \omega_p(b\beta) \right| = \delta_{\alpha,\beta} \delta_{a,b} + \frac{1}{\sqrt{d}}(1 - \delta_{a,b}) \quad (15)$$

to be compared with (2). Then, the  $|a\alpha\rangle$  vectors can be constructed once the  $\mathbf{w}(a\alpha)$  vectors are known. The solution, in matrix form, is

$$|a\alpha\rangle = \begin{pmatrix} \omega_0(a\alpha) \\ \omega_1(a\alpha) \\ \vdots \\ \omega_{d-1}(a\alpha) \end{pmatrix} \quad (16)$$

$$a = 0, 1, \dots, d \quad \alpha = 0, 1, \dots, d - 1 \quad (17)$$

Therefore, we can construct a complete set  $\{B_a : a = 0, 1, \dots, d\}$  of  $d + 1$  MUBs from the knowledge of  $d(d + 1)$  vectors  $\mathbf{w}(a\alpha)$ . Note that, for fixed  $a$  and  $\alpha$ , the  $|a\alpha\rangle$  vector is an eigenvector of the  $M_{a\alpha}$  matrix with the eigenvalue 1. This establishes a link with the above-mentioned diagonalization procedure.

### 4. A PARALLEL PROBLEM

The present work takes its origin in [6] where some similar developments were achieved for the search of a SIC POVM. Symmetric informationally complete positive-operator-valued measures play an important role in quantum information. Their existence in arbitrary dimension is still the object of numerous studies (see for instance [7]).

A SIC POVM in dimension  $d$  can be defined as a set of  $d^2$  nonnegative operators  $P_x = |\Phi_x\rangle\langle\Phi_x|$  acting on  $\mathbb{C}^d$  and satisfying

$$\frac{1}{d} \sum_{x=1}^{d^2} P_x = I \quad (18)$$

and

$$\text{Tr}(P_x P_y) = \frac{d\delta_{x,y} + 1}{d + 1} \quad (19)$$

where  $I$  is the identity operator. The search for such a SIC POVM amounts to find  $d^2$  vectors  $|\Phi_x\rangle$  in  $\mathbb{C}^d$  satisfying

$$\frac{1}{d} \sum_{x=1}^{d^2} |\Phi_x\rangle\langle\Phi_x| = I \quad (20)$$

and

$$|\langle\Phi_x|\Phi_y\rangle| = \sqrt{\frac{d\delta_{x,y} + 1}{d + 1}} \quad (21)$$

with  $x, y = 1, 2, \dots, d^2$ . The  $P_x$  operator can be developed as

$$P_x = \sum_{p=0}^{d-1} \sum_{q=0}^{d-1} v_{pq}(x) E_{pq} \quad (22)$$

so that the determination of  $d^2$  operators  $P_x$  (or  $d^2$  vectors  $|\Phi_x\rangle$ ) is equivalent to the determination of  $d^2$  vectors  $\mathbf{v}(x)$ , of components  $v_{pq}(x)$ , in  $\mathbb{C}^{d^2}$ . In the spirit of the preceding sections, we have the following result.

**Proposition 3.** For  $d \geq 2$ , finding a SIC POVM in  $\mathbb{C}^d$  (if it exists) is equivalent to finding  $d^2$  vectors  $\mathbf{v}(x)$  in  $\mathbb{C}^{d^2}$ , of components  $v_{pq}(x)$  such that

$$\frac{1}{d} \sum_{x=1}^{d^2} v_{pq}(x) = \delta_{p,q}, \quad p, q \in \mathbb{Z}/d\mathbb{Z} \quad (23)$$

$$\sum_{p=0}^{d-1} \sum_{q=0}^{d-1} \overline{v_{pq}(x)} v_{pq}(y) = \frac{d\delta_{x,y} + 1}{d+1} \quad (24)$$

and

$$v_{pq}(x) = \nu_p(x) \overline{\nu_q(x)}, \quad p, q \in \mathbb{Z}/d\mathbb{Z} \quad (25)$$

where  $x, y = 1, 2, \dots, d^2$ .

## 5. CONCLUDING REMARKS

The equivalence discussed in this work of the two ways of looking at MUBs amounts in some sense to the equivalence between the search for equiangular lines in  $\mathbb{C}^d$  and for equiangular vectors in  $\mathbb{C}^{d^2}$  (cf. [8]). Equiangular lines in  $\mathbb{C}^d$  correspond to

$$|\langle a\alpha | b\beta \rangle| = \frac{1}{\sqrt{d}} \text{ for } a \neq b \quad (26)$$

while equiangular vectors in  $\mathbb{C}^{d^2}$  correspond to

$$\mathbf{w}(a\alpha) \cdot \mathbf{w}(b\beta) = \frac{1}{d} \text{ for } a \neq b \quad (27)$$

where the  $\mathbf{w}(a\alpha) \cdot \mathbf{w}(b\beta)$  inner product in  $\mathbb{C}^{d^2}$  is defined as

$$\mathbf{w}(a\alpha) \cdot \mathbf{w}(b\beta) = \sum_{p=0}^{d-1} \sum_{q=0}^{d-1} \overline{w_{pq}(a\alpha)} w_{pq}(b\beta) \quad (28)$$

Observe that the modulus disappears and the  $1/\sqrt{d}$  factor is replaced by  $1/d$  when passing from (26) to (27). It was questioned in [5] if the equiangular vectors approach can shed light on the still unsolved question to know if one can find  $d+1$  MUBs when  $d$  is not a (strictly positive) power of a prime integer. In the case where  $d$  is not a power of a prime, the impossibility of finding  $d(d+1)$  vectors  $\mathbf{w}(a\alpha)$  or  $d(d+1)$  matrices  $M_{a\alpha}$  satisfying the conditions in Propositions 1 and 2 would mean that  $d+1$  MUBs do not exist in  $\mathbb{C}^d$ . However, it is hard to know if one approach is better than the other. It is the hope of the author that the equiangular vectors approach be tested in the  $d=6$  case for which one knows only three MUBs instead of  $d+1=7$  in spite of numerous numerical studies (see [9–11] and references therein for an extensive list of related works).

Similar remarks apply to SIC POVMs. The existence problem of SIC POVMs in arbitrary dimension is still unsolved although SIC POVMs have been constructed in every dimension  $d \leq 67$  (see [7] and references therein). For SIC POVMs, the equiangular lines in  $\mathbb{C}^d$  correspond to

$$|\langle \Phi_x | \Phi_y \rangle| = \frac{1}{\sqrt{d+1}} \text{ for } x \neq y \quad (29)$$

and the equiangular vectors in  $\mathbb{C}^{d^2}$  to

$$\mathbf{v}(x) \cdot \mathbf{v}(y) = \frac{1}{d+1} \text{ for } x \neq y \quad (30)$$

where the  $\mathbf{v}(x) \cdot \mathbf{v}(y)$  inner product in  $\mathbb{C}^{d^2}$  is defined as

$$\mathbf{v}(x) \cdot \mathbf{v}(y) = \sum_{p=0}^{d-1} \sum_{q=0}^{d-1} \overline{v_{pq}(x)} v_{pq}(y) \quad (31)$$

The parallel between MUBs and SIC POVM characterized by the couples of equations (26)–(29), (27)–(30) and (28)–(31) should be noted. These matters shall be the subject of a future work.

## ACKNOWLEDGEMENTS

The material contained in the present note was planned to be presented at the eleventh edition of the workshop *Analytic and Algebraic Methods in Physics* (AAMP XI). Unfortunately, the author was unable to participate in AAMP XI. He is greatly indebted to Miloslav Znojil for suggesting him to submit this work to Acta Polytechnica.

## REFERENCES

- [1] A. Vourdas. Quantum systems with finite Hilbert space. *Rep Prog Phys* **67**(3):267–320, 2004.
- [2] J. Tolar, G. Chadzitaskos. Feynman’s path integral and mutually unbiased bases. *J Phys A: Math Theor* **42**(24):1–11, 2009.
- [3] M. R. Kibler. An angular momentum approach to quadratic Fourier transform, Hadamard matrices, Gauss sums, mutually unbiased bases, unitary group and Pauli group. *J Phys A: Math Theor* **42**(35):1–28, 2009.
- [4] T. Durt, B.-G. Englert, I. Bengtsson, K. Życzkowski. On mutually unbiased bases. *Int J Quantum Inf* **8**(4):535–640, 2010.
- [5] M. R. Kibler. Equiangular vectors approach to mutually unbiased bases. *Entropy* **15**(5):1726–1737, 2013.
- [6] O. Albouy, M. R. Kibler. A unified approach to SIC-POVMs and MUBs. *J Russian Laser Res* **28**(5):429–438, 2007.
- [7] D. M. Appleby, C. A. Fuchs, H. Zhu. Group theoretic, Lie algebraic and Jordan algebraic formulations of the SIC existence problem. [arXiv:1312.0555v1](https://arxiv.org/abs/1312.0555v1).
- [8] C. Godsil, A. Roy. Equiangular lines, mutually unbiased bases, and spin models. *Eur J Combin* **30**(1):246–262, 2009.
- [9] P. Butterley, W. Hall. Numerical evidence for the maximum number of mutually unbiased bases in dimension six. *Phys Lett A* **369**(1-2):5–8, 2007.
- [10] S. Brierley, S. Weigert. Constructing mutually unbiased bases in dimension six. *Phys Rev A* **79**(5):1–13, 2009.
- [11] D. McNulty, S. Weigert. On the impossibility to extend triples of mutually unbiased product bases in dimension six. *Int J Quantum Inf* **10**(5):1–11, 2012.